

Chap 23 Electric fields
Chapter 22. Electric Charges

23-1

Electromagnetism { Electric (Electricity)
E + M { magnetic (Magnetism)

Electricity magnetism

↓
1820

Oersted : moving current create magnetic field

↓
Michael Faraday

James C. Maxwell (put in Mathematic form)

↓
Maxwell's Equations (Table 32-1, P. 803)

1. Electric Charges . — intrinsic characteristic of the fundamental charges

- Static Cling
- Spark between fingers and door knobs in a dry day
- lightning in the sky

2. positive charges — assigned arbitrarily by J. Franklin
negative charges —
electric neutral — balance between ~~positive~~ + negative charges
no net charges.

3. Interaction of Charged objects through electric force.

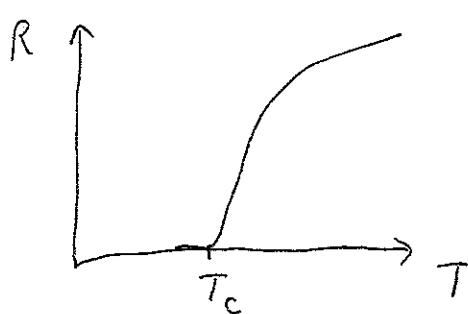
- Charges with the same electrical sign repel each other.
— Charges with the opposite sign attract each other.
- Coulomb's law, (electrostatic force , electric force)

4. Conductor - negative charge move freely

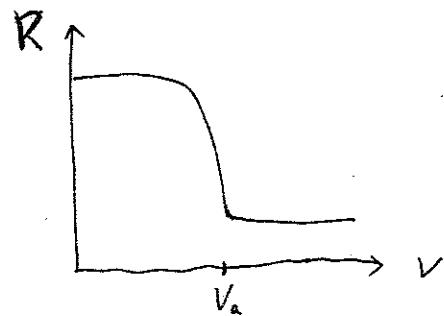
insulator - charges don't move

Semiconductor - at some apply voltage Charge will move
Si. Ge. ...

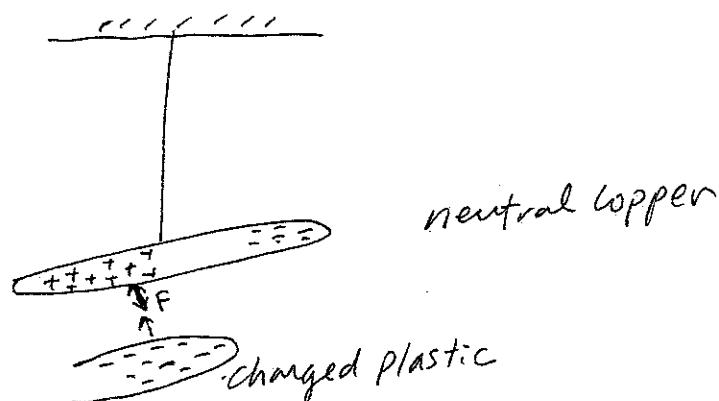
Superconductor -



Superconductor

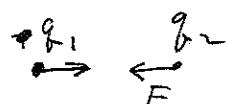


Semiconductor



5. Coulomb's law - Charles Augustin Coulomb,

$$F = k \frac{q_1 q_2}{r^2}$$



$$\rightarrow F = k \frac{q_1 q_2}{r^2}$$

Compared with
Newton's force

$$F = G \frac{m_1 m_2}{r^2}$$

$$\rightarrow F = \pm k \frac{q_1 q_2}{r^2}$$

"+" repel
"-" attract

Compare
All forces

k = electrostatic constant

$$= \frac{1}{4\pi\epsilon_0} = 8.99 \times 10^9 \text{ N} \cdot \frac{\text{m}^2}{\text{C}^2}$$

$$\epsilon_0 = 8.85 \times 10^{-12} \text{ C}^2/\text{N} \cdot \text{m}^2 = \text{permittivity constant.}$$

C : in Coulomb. SI unit.

9: SI unit, Coulomb

23-3

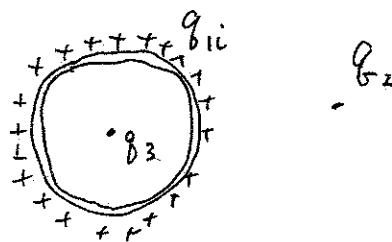
1 coulomb = $1 \text{ A} - \boxed{1 \text{ C}}$ → 1 A in one second.

$$dq = i dt$$

$$\rightarrow F_{12} = k \frac{q_1 q_2}{r^2} = \frac{1}{4\pi\epsilon_0} \frac{q_1 q_2}{r^2}$$

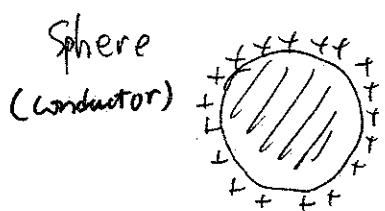
Electric force follow the same vector sum as that of gravitational force, ie, $\vec{F}_1 = \vec{F}_{12} + \vec{F}_{13} + \vec{F}_{14} + \dots$, and the force ~~act~~ on one particle is the sum of all other particles force on this one.

6. Shell:



the force on q_2 is the sum of all q_i 's as if they were centered on the center of the shell

But q_3 will exert no force due to all q_i 's in the shell.



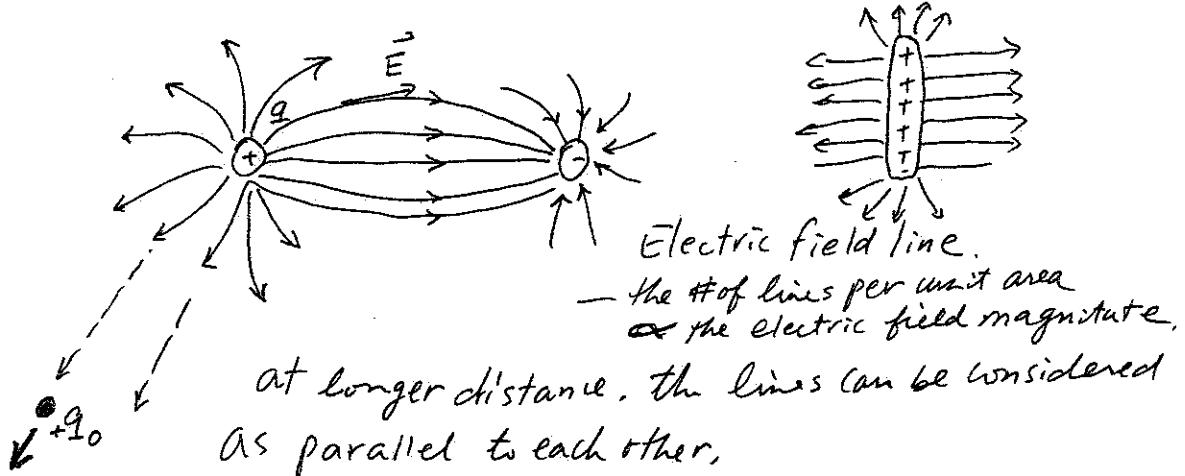
Charges will spread over the surface of a sphere uniformly because charges will repel one another

7 Charge is "Quantized" — ie, $q = ne$, e = charge of an electron
 e is the elementary charge.

$$= 1.6 \times 10^{-19} \text{ Coulombs}$$

8 Charge is conserved

Electric field — Set up by charge, within it, another charge is subject to force according to Coulomb's law



If we place a test charge q_0 .

$$F = k \frac{q q_0}{r^2}$$

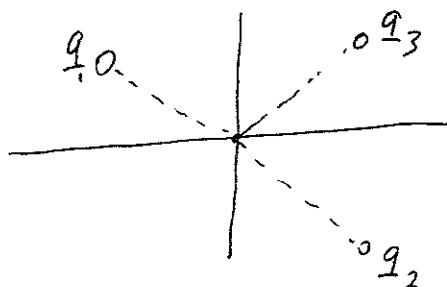
$$\vec{E} = \frac{\vec{F}}{q_0} = k \frac{q}{r^2}, \text{ Electric field. } \left[\frac{N}{C} \right]$$

- ① Electric field lines are drawn from positive charge to negative charge
- ② The # density of lines is proportional to the magnitude of the electric field.

1. Electric field — discrete charge

① Point charge $\Rightarrow E = \frac{F}{q} = \frac{1}{4\pi\epsilon_0} \frac{|q|}{r^2}$

② Multiple point charges



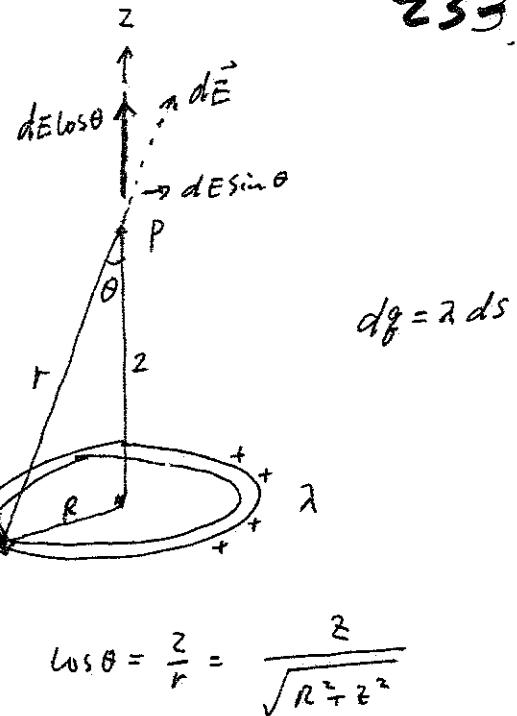
\vec{E}_0 at the origin
 $\vec{E}_0 = \vec{E}_1 + \vec{E}_2 + \vec{E}_3$

Charge distribution

23-5

④ Ring of charge distribution.

$$d\vec{E} = \frac{1}{4\pi\epsilon_0} \frac{dq}{r^2} = \frac{1}{4\pi\epsilon_0} \frac{\lambda ds}{(z^2 + R^2)}$$



$$\begin{aligned}\vec{E} &= \int d\vec{E} \cos \theta = \int \frac{1}{4\pi\epsilon_0} \frac{\lambda ds}{(z^2 + R^2)} \cdot \frac{z}{\sqrt{z^2 + R^2}} ds \\ &= \frac{\lambda}{4\pi\epsilon_0} \frac{z}{(z^2 + R^2)^{3/2}} \cdot \int_0^{2\pi R} ds \\ &= \frac{2\pi R \lambda}{4\pi\epsilon_0} \frac{z}{(z^2 + R^2)^{3/2}} = \frac{1}{4\pi\epsilon_0} \frac{q z}{(z^2 + R^2)^{3/2}}.\end{aligned}$$

$$\text{if } z \gg R, \quad z^2 + R^2 \approx z^2$$

$$\vec{E} = \frac{1}{4\pi\epsilon_0} \frac{q}{z^2}$$

At a large distance, a ring distribution can be approximated as the point charge.

Note: i) Always check the answer at extreme case

ii) Do Sample problem 23-5 for a section of ring distribution

⑤ Charged disk (charge on its upper surface)

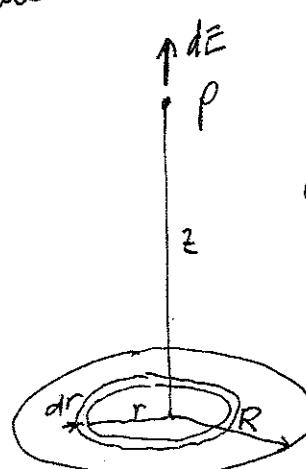
$$dq = \sigma dA = \sigma (2\pi r dr)$$

$$d\vec{E} = \frac{z \sigma 2\pi r dr}{4\pi\epsilon_0 (z^2 + r^2)^{3/2}}$$

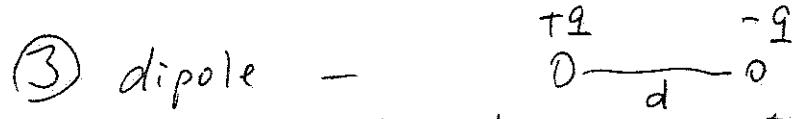
$$= \frac{\sigma z}{4\epsilon_0} \frac{2r dr}{(z^2 + r^2)^{3/2}}$$

$$E = \int_{r=0}^{r=R} d\vec{E} = \frac{\sigma z}{4\epsilon_0} \int_0^R \frac{2r dr}{(z^2 + r^2)^{3/2}}$$

$$= \frac{\sigma z}{4\epsilon_0} \left[\frac{(z^2 + r^2)^{-1/2}}{-1} \right]_0^R = \frac{\sigma}{2\epsilon_0} \left(1 - \frac{z}{\sqrt{z^2 + R^2}} \right)$$



σ = Charge density



two charges separated by a distance, have opposite sign are called electric dipole

$$\vec{E} = \vec{E}_{(+)} + \vec{E}_{(-)} = \text{electric field @ pt.P.}$$

$$= \frac{1}{4\pi\epsilon_0} \left[\frac{q}{r_+^2} - \frac{q}{r_-^2} \right]$$

$$= \frac{q}{4\pi\epsilon_0} \left[\frac{1}{(z - \frac{d}{2})^2} - \frac{1}{(z + \frac{d}{2})^2} \right]$$

$$= \frac{q}{4\pi\epsilon_0 z^2} \left[\left(1 - \frac{d}{2z}\right)^{-2} - \left(1 + \frac{d}{2z}\right)^{-2} \right], \quad z \gg d, \quad \frac{d}{2z} \ll 1$$

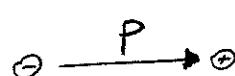
$$= \frac{q}{4\pi\epsilon_0 z^2} \left[\left(1 + \frac{2d}{2z}(1) + \dots\right) - \left(1 - \frac{2d}{2z}(1) + \dots\right) \right]$$

$$\approx \frac{q}{4\pi\epsilon_0 z^2} \cdot \frac{2d}{3}$$

$$= \frac{1}{2\pi\epsilon_0} \frac{qd}{z^3}$$

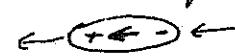
$$= \frac{1}{2\pi\epsilon_0} \frac{\vec{P}}{z^3}$$

$P \equiv qd = \text{electric dipole moment.}$



in general

i) $\vec{E} \propto \frac{1}{r^3}$ for all dipole at all distance

ii) The direction of the dipole is taken from the negative to the positive charge 

iii) At a distance, \vec{E} is the same direction of \vec{P}

iv) $E \propto \frac{1}{r^2}$ for a point charge.

v) $H_2O. \quad p = 6.2 \times 10^{-30} \text{ C.m}$



23-7

for a charged disk $E = \frac{\sigma}{2\epsilon_0} \left(1 - \frac{z}{\sqrt{z^2 + R^2}} \right)$

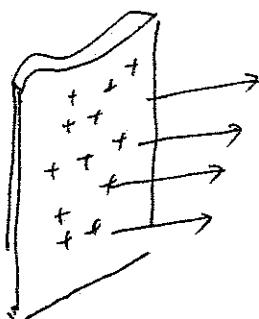
⑥ Infinite large Sheet

from E for the disk, $E = \frac{\sigma}{2\epsilon_0} \left(1 - \frac{z}{\sqrt{z^2 + R^2}} \right)$

if $R \gg z$, $R \rightarrow \infty$, $\frac{z}{\sqrt{z^2 + R^2}} \approx 0$

$E \rightarrow \frac{\sigma}{2\epsilon_0}$ — Electric field for a large sheet.

$E \propto \sigma$, from different pt of view.

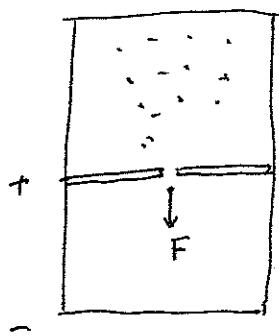


\vec{E} must be uniform and $\propto \sigma$, the Surface charge density -

2. Electric force — for a pt charge in an electric field

$$\vec{F} = q\vec{E} \quad \text{Note: pt charge doesn't mean it is Unit Charge.}$$

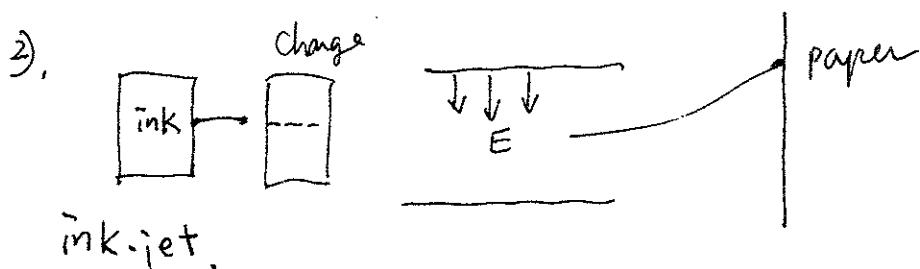
Millikan Oil drop exp.



$$F = qE$$

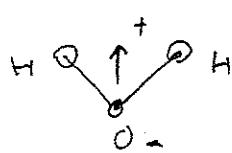
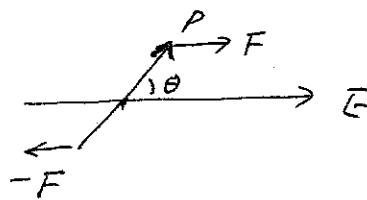
$$q = ne$$

$$e = 1.6 \times 10^{-19} \text{ Coulomb.}$$



ink-jet.

3. dipole in an electric field.

 H_2O has permanent dipole.

Force due to the interaction of electric field and charge in a dipole, the ^{net} force is zero. but exert a net torque about its center of mass.

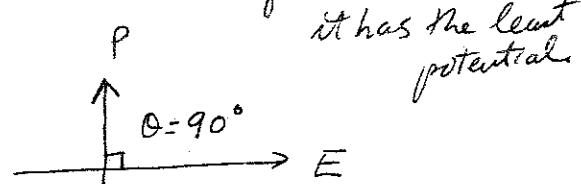
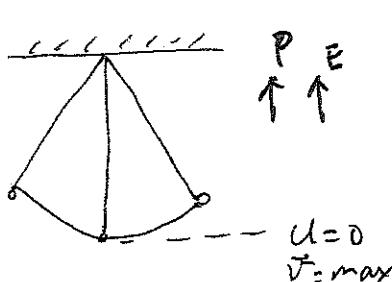
$$\tau = F \frac{d}{2} \sin\theta + F \frac{d}{2} \sin\theta = F d \sin\theta$$

$$= qE d \sin\theta$$

$$\tau = pE \sin\theta = \vec{P} \times \vec{E} = -pE \sin\theta \quad (\text{negative torque})$$

→ this give rise to the rotation of the molecules.

But the potential energy is minimum when the dipole is lined up along the external field. i.e. $\tau = p \times E = 0$, $U = 0$ compared to a pendulum, at this point, the dipole is at its equilibrium position



$U = \infty \text{ max}$ defined potential = $\tau = \infty$ max. define the potential.

$$\therefore U = -W = - \int_{90^\circ}^0 \tau d\theta = + \int_{90^\circ}^0 pE \sin\theta d\theta - \text{potential at } \theta = 90^\circ$$

$$\therefore U = -p \cdot E \cos\theta = -p \cdot E \left[\cos 90^\circ \right] = -p \cdot E \cos 90^\circ$$

4. Micro waver.

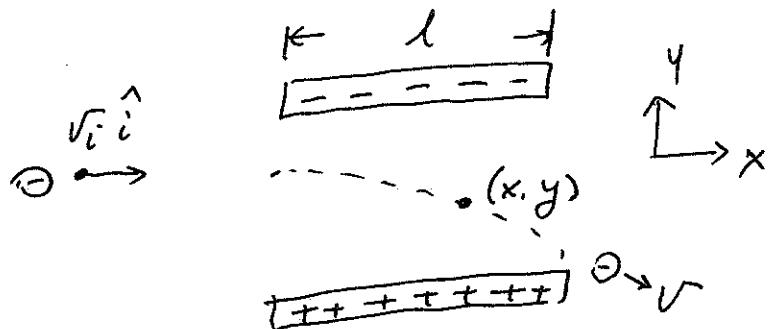
Note: $\tau = -pE \sin\theta$
negative torque

as in the figure,
it gives rise to a clockwise
rotation

23.7

Motion of a charge particle in a Uniform electric field

$$F_E = qE = ma \rightarrow a_y = \frac{qE}{m}$$



$$v_x = v_i = \text{constant}$$

$$v_y = a_y t = -\frac{eE}{mc} t$$

$$x_f = v_i t$$

$$y_f = \frac{1}{2} a_y t^2 = -\frac{1}{2} \frac{eE}{mc} t^2$$

23.5. Electric field of a continuous Charge distribution

$$\vec{\Delta E} = k_e \frac{\Delta q}{r^2} \hat{r}$$

$$\therefore \vec{E} \sim k_e \sum_i \frac{\Delta q_i}{r_i^2} \hat{r}_i$$



if the charge is continuous distributed

$$\begin{aligned}\vec{E} &= k_e \lim_{\Delta q_i \rightarrow 0} \sum_i \frac{\Delta q_i}{r_i^2} \hat{r}_i \\ &= k_e \int \frac{dq}{r^2} \hat{r}\end{aligned}$$