

Dividing by  $\frac{dI}{dt}$ ,  $\mathcal{E} = -M_{21} \frac{dI_1}{dt}$

We could get the definition of  $M$   $M \equiv \frac{\mu_0}{4\pi} \int_{\lambda} \int_{\lambda'} \frac{d\vec{s} \cdot d\vec{\lambda}'}{r}$   
 where  $d\vec{s}$  &  $d\vec{\lambda}'$  are two element of length (Eq 7.2.3)  
 &  $r$  is the distance between them. eg 7.22

This is known as Neumann's formula.  $M_{12} = \frac{\mu_0}{4\pi} \int_{\lambda} \int_{\lambda'} \frac{d\vec{\lambda} \cdot d\vec{s}'}{r}$

6/3 (E)

$\Rightarrow M_{12} = M_{21}$

\* Final Exam. 6/17 (E) 10:00 AM.

$\Phi = MI$   
 $\downarrow$  flux  $\downarrow$  mutual inductance

§ 7.2.4 Energy in magnetic field  $\Rightarrow$  Energy + Density of Energy

To obtain the energy of an electrostatic field, we calculate the work done by  $\vec{E}$  in moving charge increment  $dq$ .

Then the increment of  $dq$  can be described as  $dq = I dt$   
 \*  $dW = -\mathcal{E} dt$ , through the circuit is  $\Rightarrow dW = -\mathcal{E} I = -L I \frac{dI}{dt}$   
 then we can integrate the time function of work.

$W = \frac{1}{2} L I^2(t)$

1. If  $dW = I d\Phi$  is equal to the change in magnetic energy.
2. If there are  $n$  turns circuit, then the work done against the induced emf is given by  $dW = \sum_{i=1}^n I_i d\Phi_i$

\* Magnetic Energy in terms of field vectors.

For simplicity, we assume that each circuit consist of a single loop, then the flux  $\Phi$

$\Phi_i = \int \vec{B} \cdot \hat{n} da = \int \nabla \times \vec{A} \cdot \hat{n} da = \int_{\text{closed loop}} \vec{A} \cdot d\vec{\lambda}$

$\Rightarrow$  input the magnetic energy.

For  $n$  turns, the function of  $dW = \sum I_i \Phi_i$  can be replaced as  $W = \frac{1}{2} I \Phi$ .  $\Rightarrow$  Magnetic energy  
 $\Rightarrow U = \frac{1}{2} \sum I_i \Phi_i \quad U = \frac{1}{2} \int I_i \vec{A}_i \cdot d\vec{l}$

\* The second step: we can change  $I_i dl_i = (\vec{J} \cdot d\vec{a}) \cdot dl$   
 $= \vec{J} \cdot d\vec{v}_i$

$$\& \oint_{\text{closed path}} \vec{A}_i = \int_{\text{Volume}} \vec{J} \cdot d\vec{v}_i \Rightarrow \text{then } U = \frac{1}{2} \int \vec{J} \cdot \vec{A} dV$$

\* Mathematical method of  $\vec{B} = \nabla \times \vec{A}$ ,  $\nabla \times \vec{B} = \mu_0 \vec{J}$

$$\text{Re-calculated } \nabla \cdot (\vec{A} \times \vec{B}) = \vec{B} \cdot (\nabla \times \vec{A}) - \vec{A} \cdot (\nabla \times \vec{B}) \\ = \vec{B} \cdot \vec{B} - \vec{A} \cdot \mu_0 \vec{J}$$

$$\Rightarrow \boxed{\vec{B} \cdot \vec{B} - \nabla \cdot (\vec{A} \times \vec{B}) = \vec{A} \cdot \mu_0 \vec{J}} \quad \text{VIP}$$

$U = \frac{1}{2} \int \vec{J} \cdot \vec{A} dV$ , we obtain the expression of

$$U = \frac{1}{2\mu_0} \int \vec{B} \cdot \vec{B} dV - \frac{1}{2\mu_0} \int \nabla \cdot (\vec{A} \times \vec{B}) dV$$

$$= \frac{1}{2\mu_0} \int \vec{B} \cdot \vec{B} dV - \frac{1}{2\mu_0} \int \vec{A} \times \vec{B} \cdot \hat{n} da$$

Explan.

The integrations on the right are to be taken over

**Uniform** the entire volume occupied by the current with the surface

Because  $\vec{B}$  falls off at fast as  $\frac{1}{r^2}$

$\vec{A}$

"

$\frac{1}{r}$

Surface  $da \sim r^2$

Then the second term of surface integral vanishes.

$$U = \frac{1}{2\mu_0} \int \vec{B} \cdot \vec{B} dV, \text{ when } \vec{B} = \mu_0 \vec{H}$$

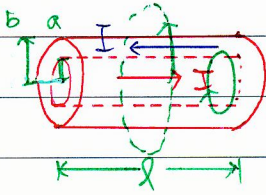
$$= \frac{1}{2} \int \vec{H} \cdot \vec{B} dV, \text{ we may define the}$$

\* energy density in magnetic field by  $U = \frac{1}{2} \vec{H} \cdot \vec{B}$

$$u = \frac{B^2}{2\mu}$$



§ 7.13 Example. A long co-axial cables carries current  $I$ .



The outer current  $I \leftarrow$

inner current  $I \rightarrow$

Find the magnetic energy stored in a section of length  $l$ .

\* According to Ampere's law.  $B = \frac{\mu_0 I}{2\pi r} \hat{\phi}$

\* The magnetic energy per unit volume

$$U = \frac{1}{2\mu_0} B^2 = \frac{\mu_0 I^2}{8\pi^2 r^2}$$

\* The unit volume of cylinder  $U = \int_{\text{volume}} dv$ .  $dv = l \cdot 2\pi r dr$

So the magnetic energy is  $U = \int U dv = \int_a^b \frac{\mu_0 I^2}{8\pi^2 r^2} \cdot 2\pi r dr \cdot l$

$$= \frac{\mu_0 I^2 l}{4\pi} \int_a^b \frac{1}{r} dr = \frac{\mu_0 I^2 l}{4\pi} \ln \frac{b}{a}, \text{ represented with } \Delta \cdot I.$$

$$U = \frac{1}{2} \Delta I^2 \quad \Delta = \frac{\mu_0 l}{2\pi} \ln \frac{b}{a}$$

HW 06/10 return

P7.18 P7.24

P7.20 P7.30

§ 7.3 Maxwell Equation.

7.31. Electrodynamics equation. Before Maxwell

A. Gauss's Law  $\nabla \cdot \vec{E} = \frac{\rho}{\epsilon_0}$

$\nabla \cdot \vec{B} = 0$

Faraday's Law  $\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$

Ampere's Law  $\nabla \times \vec{B} = \mu_0 \vec{J}$

These equations represent the state of EM theory over a century.

The old eqs. with the old rule that divergency of curl is always zero

$$(1) \nabla \cdot (\nabla \times \vec{E}) = \nabla \cdot \left( -\frac{\partial \vec{B}}{\partial t} \right) = -\frac{\partial}{\partial t} (\nabla \cdot \vec{B}) = 0$$

$$(2) \nabla \cdot (\nabla \times \vec{B}) = (\nabla \cdot \vec{J}) \mu_0 = 0 \quad ?$$

Zero only at static steady state; Non zero at dynamics state.



Maxwell equation & Electromagnetic wave in vacuum.

1. Gauss's Law  $\nabla \cdot \vec{E} = \frac{\rho}{\epsilon_0}$
2. Absence of magnetic monopole  $\nabla \cdot \vec{B} = 0$
3. Ampere's Law  $\nabla \times \vec{B} = \mu_0 \vec{J}$
4. Faraday's Law  $\nabla \times \vec{E} + \frac{\partial \vec{B}}{\partial t} = 0$
5. The continuity eq.  $\frac{\partial \rho}{\partial t} + \nabla \cdot \vec{J} = 0$

6/4(IV) 6/17 (E) Final Exam of EM.

§ 7.3.1

(1) Ampere's Law:  $\nabla \times \vec{B} = \mu_0 \vec{J} \Rightarrow \nabla \cdot (\nabla \times \vec{B}) = \mu_0 (\nabla \cdot \vec{J}) = 0$   
 at steady, valid for dynamics?

(2) Continuity equation:  $\nabla \cdot \vec{J} + \frac{\partial \rho}{\partial t} = 0$ . Charge increment  $\neq 0$

§ 7.3.2

How Maxwell fixed Ampere's Law for Electrodynamics?

<1> Continuity equation & Ampere's Law, they were all valid even in time varying situation & realize Ampere's Law was inconsistent with the continuity equation.

$\Leftrightarrow \nabla \cdot (\nabla \times \vec{B}) = 0, \text{ \& } \nabla \cdot \vec{J} \neq 0$

This indeed true if  $\rho$  does not change with time. But it is not true when  $\rho$  is changing with time.

$\frac{\Delta \rho}{\Delta t} = \text{constant}$  velocity

$\frac{\rho}{t} = \text{constant}$  Speed

$\frac{\Delta \rho + \rho}{\Delta t} \neq \text{constant}$

Accel



<3> Using Gauss's Law, we can rewrite the continuity eq.

$$\text{as } \nabla \cdot \vec{J} = -\frac{\partial \rho}{\partial t} = -\frac{\partial}{\partial t} [\nabla \cdot \epsilon_0 \vec{E}]$$

$$\text{Replace } \nabla \cdot [\vec{J} + \frac{\partial}{\partial t} \epsilon_0 \vec{E}] = 0 \Rightarrow \nabla \cdot \vec{J}_{\text{Maxwell}} = 0.$$

<4> So the divergence of curl B can be rewritten as

$$\nabla \cdot (\nabla \times \vec{B}) = \nabla \cdot (\mu_0 \vec{J}_{\text{Maxwell}}) = \nabla \cdot [\mu_0 \vec{J} + \frac{\partial}{\partial t} \mu_0 \epsilon_0 \vec{E}] = 0$$

Maxwell saw that if Ampere's Law is modified by the addition of a new term the time derivative term.

$$\Rightarrow \nabla \times \vec{B} = \mu_0 \vec{J} + \frac{\partial}{\partial t} (\underbrace{\mu_0 \epsilon_0}_{\text{Mag.}} \underbrace{\vec{E}}_{\text{Ele.}})$$

(1) Is valid for Steady-state phenomena is also compatible with the equation of continuity of time-dependent fields (dynamics)

(2) The term  $\epsilon_0 \frac{\partial \vec{E}}{\partial t}$  has the dimensions of current density

$$\nabla \cdot (\epsilon_0 \vec{E} + \vec{P}) = \rho_f \Rightarrow \nabla \cdot \vec{D} = \rho_f$$

§ Have a look in Gauss's Law for Electric & Magnetic fields.

$$(1) \nabla \cdot \vec{E} = \frac{\rho}{\epsilon_0} + (2) \nabla \cdot \vec{B} = 0 \quad (\mu_0 \rho_m = 0)$$

Using the free space ( $e$  or  $m = 0$ )

$$\nabla \cdot \vec{E} = 0 \quad \nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

$$\nabla \cdot \vec{B} = 0 \quad + \quad \nabla \times \vec{B} = \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t}$$

$$\text{replace } \vec{E} \text{ by } \vec{B} \rightarrow \Rightarrow \nabla \cdot \vec{B} = 0, \quad \nabla \times \vec{B} = \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t}$$

$$\text{replace } \vec{B} \text{ by } -\mu_0 \epsilon_0 \vec{E} \rightarrow \Rightarrow \nabla \cdot \vec{E} = 0, \quad \nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

There are something missing from  $\nabla \cdot \vec{B} = 0$  &  $\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$

if we had  $\rho_e, \rho_m$  &  $m$ . (In non-free space)

$$\text{We have } \textcircled{1} \nabla \cdot \vec{E} = \frac{\rho_e}{\epsilon_0} \quad \textcircled{3} \nabla \times \vec{E} = -\mu_0 \vec{J}_m - \frac{\partial \vec{B}}{\partial t}$$

$$\textcircled{2} \nabla \cdot \vec{B} = \mu_0 \rho_m \quad \textcircled{4} \nabla \times \vec{B} = \mu_0 \vec{J}_e + \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t}$$

$\rho_m$ : magnetic charge density.

$\vec{J}_m$ : magnetic current density.



"Maxwell's equation beg for the existance of magnetic charge" 7.35, 7.36