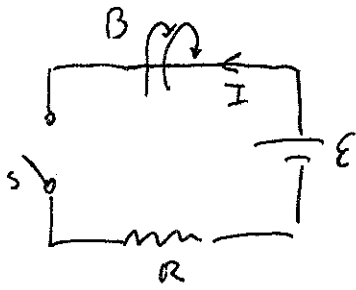


32.1 Self inductance

The current does not go from 0 to max, immediately



① when the switch is off, the current increases gradually, which causes an increase in the magnetic field

② the increased magnetic flux creates an induced emf opposite the change of the original magnetic flux.

→ also ~~the~~ gradually increased current

③ the effect is called self-induction, creating an self-induced emf.

$$\mathcal{E}_L \propto \frac{dI}{dt}$$

$$\mathcal{E}_L = -L \frac{dI}{dt} \rightarrow \text{the self-induced emf is proportional to the time rate change of the current}$$

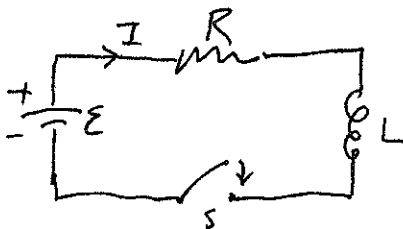
$$L \equiv \text{inductance} \equiv \text{---}$$

But from Faraday's law  $\mathcal{E}_L = -N \frac{d\Phi_B}{dt} = -L \frac{dI}{dt}$

$$\therefore L = \frac{N\Phi_B}{I} \quad \text{or} \quad L = -\frac{\mathcal{E}_L}{\frac{dI}{dt}} \quad [\text{Henry}] = 1 \frac{\text{V}\cdot\text{s}}{\text{A}}$$

32.2 RL circuit

An inductor in a circuit opposes changes in the current in that circuit.



Using Kirchhoff's law

$$\mathcal{E} - IR - L \frac{dI}{dt} = 0$$

$$\text{let } x = \frac{\mathcal{E}}{R} - I \rightarrow dx = -dI$$

$$\therefore x + \frac{L}{R} \frac{dx}{dt} = 0$$

$$\frac{dx}{x} = -\frac{R}{L} dt$$

$$x + \frac{L}{R} \frac{dx}{dt} = 0$$

$$\frac{dx}{x} = -\frac{R}{L} dt$$

$$\int_{x_0}^x \frac{dx}{x} = -\frac{R}{L} \int dt \rightarrow \ln \frac{x}{x_0} = -\frac{R}{L} t$$

$$\rightarrow x = x_0 e^{-\frac{R}{L} t}$$

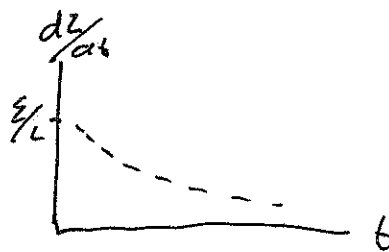
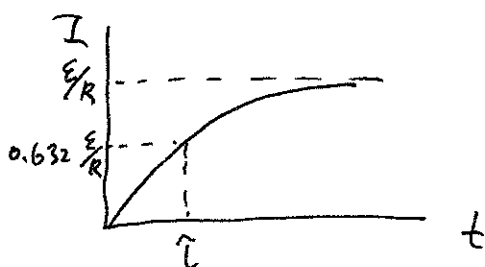
$$I = 0 \text{ at } x = 0$$

$$x_0 = \frac{\mathcal{E}}{R}$$

$$\therefore \frac{\mathcal{E}}{R} - I = \frac{\mathcal{E}}{R} e^{-\frac{R}{L} t}$$

$$\rightarrow I = \frac{\mathcal{E}}{R} (1 - e^{-\frac{R}{L} t})$$

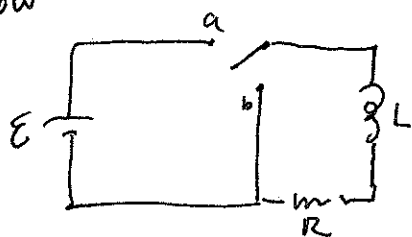
$$\text{or } I(t) = \frac{\mathcal{E}}{R} (1 - e^{-\frac{t}{\tau}}) \quad , \quad \frac{dI(t)}{dt} = \frac{\mathcal{E}}{L} e^{-\frac{t}{\tau}}$$



Switch open  $t < 0$

Switch closed  $t > 0$

Now

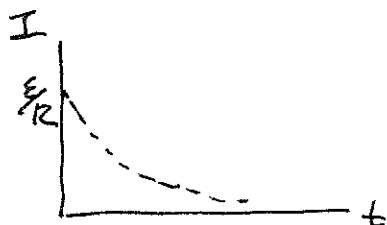


if the switch is thrown to b

$$\rightarrow \mathcal{E} = 0$$

$$\rightarrow \mathcal{E} + R + L \frac{dI}{dt} = 0$$

$$\text{then } I(t) = \frac{\mathcal{E}}{R} e^{-\frac{t}{\tau}} = I_0 e^{-\frac{t}{\tau}}$$



### 32.3 Energy in a Magnetic field

— Because the emf induced in an inductor prevents a battery from establishing an instantaneous current, the battery must provide more energy than a circuit without an inductor

$$(32.6) \quad \mathcal{E} - IR - L \frac{dI}{dt} = 0$$

$$I\mathcal{E} = I^2R + \underbrace{LI \frac{dI}{dt}}_{\text{energy rate stored in the inductor}}$$

$\underbrace{I\mathcal{E}}_{\text{energy supplied by the battery}} = \underbrace{I^2R}_{\text{rate energy is delivered to the resistor}} + LI \frac{dI}{dt}$

$$\therefore \frac{dU}{dt} = LI \frac{dI}{dt} \quad U \equiv \text{energy stored in an inductor}$$

$$\rightarrow U = \int dU = \int_0^I LI dI = L \int_0^I I dI = \frac{1}{2} LI^2$$

$$\therefore \underline{U = \frac{1}{2} LI^2} \quad \text{energy stored in an inductor}$$

Note  $U = \frac{1}{2} C (\Delta V)^2$  — energy stored in a capacitor

Consider a solenoid of  $N$  turns

$$L = \mu_0 n^2 AL \quad (32.5) \text{ example 32.1}$$

$$B = \mu_0 n I \quad (30.7)$$

$$\therefore U = \frac{1}{2} LI^2 = \frac{1}{2} \mu_0 n^2 AL \cdot \left(\frac{B}{\mu_0 n}\right)^2 = \frac{B^2}{2\mu_0} AL$$

$$u_B = \frac{U}{V} = \frac{1}{2} \frac{B^2}{\mu_0} \quad \text{in an inductor}$$

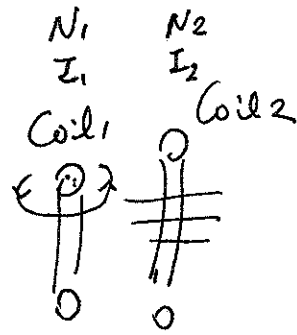
$$u_E = \frac{1}{2} \epsilon_0 E^2 \quad \text{in a capacitor}$$

32.4 Mutual Inductance

The nearby circuits create magnetic flux passing each other

→ ~~Self~~ Mutual Induction

→ Mutual Inductance Two adjacent coils



$$M_{12} \equiv N_2 \frac{\Phi_{12}}{I_1} \quad \text{Coil 2 with respect to Coil 1 due to current } I_1 \text{ at Coil 2.}$$

If  $I_1$  varies with time

The emf induced by Coil 1 in Coil 2

$$\mathcal{E}_2 = -N_2 \frac{d\Phi_{12}}{dt} = -N_2 \frac{d}{dt} \left( \frac{M_{12} I_1}{N_2} \right) = -M_{12} \frac{dI_1}{dt}$$

Similarly

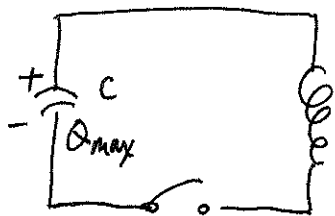
$$\mathcal{E}_1 = -M_{21} \frac{dI_2}{dt}$$

But  $M_{12} = M_{21}$

$$\therefore \mathcal{E}_2 = -M \frac{dI_1}{dt}$$

$$\mathcal{E}_1 = -M \frac{dI_2}{dt}$$

32.5 Oscillations in an LC circuit



$$U = U_C + U_L = \frac{1}{2C} Q^2 + \frac{1}{2} L I^2$$

$$\frac{dU}{dt} = \frac{d}{dt} \left( \frac{Q^2}{2C} + \frac{L I^2}{2} \right)$$

$$= \frac{Q}{C} \frac{dQ}{dt} + L I \frac{dI}{dt}$$

The total energy of the system is constant

$$\frac{dU}{dt} = \frac{Q}{C} \frac{dQ}{dt} + L I \frac{dI}{dt} = 0$$

$$\frac{Q}{C} \frac{dQ}{dt} + \frac{dQ}{dt} L \frac{d^2Q}{dt^2} = 0$$

$$\rightarrow \frac{Q}{C} + L \frac{d^2Q}{dt^2} = 0$$

$$\boxed{\frac{d^2Q}{dt^2} = -\frac{1}{LC} Q}$$

Compare  $\frac{dx^2}{dt^2} = -\frac{k}{m} x = -\omega^2 x$   
in a mechanical system

$$\rightarrow Q = Q_{\max} \cos(\omega t + \phi), \quad \omega = \frac{1}{\sqrt{LC}}$$

$$I = \frac{dQ}{dt} = -\omega Q_{\max} \sin(\omega t + \phi)$$

$$I = 0 \text{ @ } t = 0$$

$$Q = Q_{\max} \text{ @ } t = 0$$

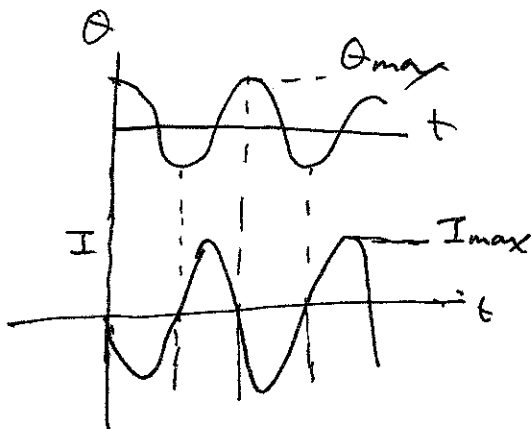
$$\therefore 0 = -\omega Q_{\max} \sin \phi \rightarrow \phi = 0$$

$$\therefore Q(t) = Q_{\max} \cos \omega t$$

$$I(t) = -\omega Q_{\max} \sin \omega t = -I_{\max} \sin \omega t$$

therefore

$$U = U_C + U_L = \frac{Q_{\max}^2}{2C} \cos^2 \omega t + \frac{1}{2} L I_{\max}^2 \sin^2 \omega t$$



$$\text{But } \frac{Q_{\max}^2}{2C} = \frac{L I_{\max}^2}{2}$$

$$\therefore U = \frac{Q_{\max}^2}{2C} (\cos^2 \omega t + \sin^2 \omega t) \\ = \frac{Q_{\max}^2}{2C}$$



32-6 The RLC circuit

①  $S_1$  closed,  $S_2$  open  
 → C charged to have  $Q_{max}$

②  $S_1$  open,  $S_2$  closed

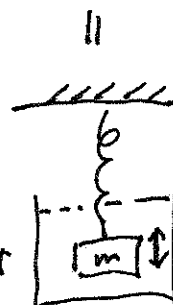
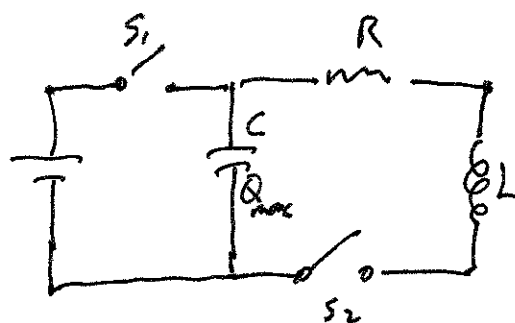
→ A ~~current~~ current establish

The total energy in C and L

is

$$U = U_C + U_L = \frac{Q^2}{2C} + \frac{1}{2} L I^2$$

But the total energy is NOT constant  
 since the resistor exists



$$-\frac{dU}{dt} = -I^2 R \quad \rightarrow \text{rate of energy transformation in a resistor}$$

$$\therefore \frac{dU_{total}}{dt} \rightarrow L I \frac{dI}{dt} + \frac{Q}{C} \frac{dQ}{dt} = -I^2 R$$

$$L I \frac{d^2 Q}{dt^2} + I^2 R + \frac{Q}{C} I = 0$$

$$L \frac{d^2 Q}{dt^2} + I R + \frac{Q}{C} = 0$$

$$L \frac{d^2 Q}{dt^2} + R \frac{dQ}{dt} + \frac{Q}{C} = 0$$

Similar to

$$m \frac{d^2 x}{dt^2} + b \frac{dx}{dt} + kx = 0 \quad (15.31)$$

— damped Oscillator

When  $R$  is small  $\rightarrow$  light damped Oscillator

P32-7

$$Q = Q_{\max} e^{-\frac{Rt}{2L}} \cos \omega_d t$$

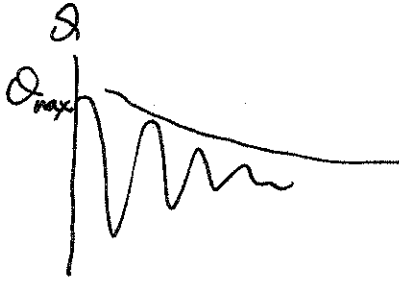
$$\omega_d = \left[ \frac{1}{LC} - \left(\frac{R}{2L}\right)^2 \right]^{\frac{1}{2}}$$

① When  $R \ll \sqrt{\frac{4L}{C}}$   $\omega_d \approx \omega$  for undamped

$$= \left[ \frac{1}{LC} \right]^{\frac{1}{2}}$$

② When  $R_c = \sqrt{\frac{4L}{C}}$  critical damped

③ When  $R > R_c$  overdamped



Check P.1021

P1022

Check Table 32.1