

1.

$$(a) \quad e \equiv \frac{W_{\text{Carnot}}}{|Q_h|} = \frac{|Q_h| - |Q_c|}{|Q_h|} = 1 - \frac{|Q_c|}{|Q_h|}$$

In a Carnot cycle, $\Delta E_{\text{in}} = 0$ (path $A \rightarrow B$)

$$\therefore |Q_h| = |-W_{AB}| \quad W_{AB} \equiv \text{work done between } A \rightarrow B$$

$$A \rightarrow B, |Q_h| = |-W_{AB}| = nR T_h \ln\left(\frac{V_B}{V_A}\right) \quad \text{absorb energy}$$

$$C \rightarrow D, |Q_c| = |-W_{CD}| = nR T_c \ln\left(\frac{V_C}{V_D}\right)$$

$$\therefore \frac{|Q_c|}{|Q_h|} = \frac{T_c}{T_h} \frac{\ln\left(\frac{V_C}{V_D}\right)}{\ln\left(\frac{V_B}{V_A}\right)} \quad \text{--- (1)}$$

$$\text{But } P_i V_i^{\gamma} = P_f V_f^{\gamma} \rightarrow T_i V_i^{\gamma-1} = T_f V_f^{\gamma-1}$$

$$\left. \begin{array}{l} B \rightarrow C \quad T_h V_B^{\gamma-1} = T_c V_C^{\gamma-1} \\ D \rightarrow A \quad T_h V_A^{\gamma-1} = T_c V_D^{\gamma-1} \end{array} \right\} \Rightarrow \left(\frac{V_B}{V_A}\right)^{\gamma-1} = \left(\frac{V_C}{V_D}\right)^{\gamma-1} \quad \therefore \frac{V_B}{V_A} = \frac{V_C}{V_D} \quad \text{--- (2)}$$

$$\text{From (1) and (2)} \quad \frac{|Q_c|}{|Q_h|} = \frac{T_c}{T_h}$$

$$\text{Therefore } e_{\text{Carnot}} = 1 - \frac{|Q_c|}{|Q_h|} = 1 - \frac{T_c}{T_h} \quad \#$$

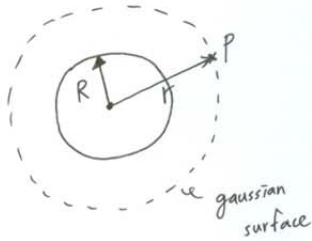
(b) For an engine to have perfect performance $e=1$

$$\Rightarrow |Q_c| = 0 \quad \text{or} \quad T_c \rightarrow 0 \quad \#$$

(c) The perfect engine requires $T_c=0$ not possible or $Q_c=0$, that means it will not give up or waste any energy. This is not possible in reality.

2.

(1) A point outside the sphere: Pick a gaussian surface bigger than the surface of the sphere



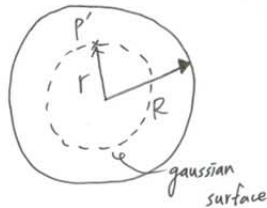
$$\Phi_E = \oint \vec{E} \cdot d\vec{A} = \frac{q_{in}}{\epsilon_0}$$

$$\Rightarrow E \oint dA = E \cdot 4\pi r^2 = \frac{Q}{\epsilon_0}$$

$$\Rightarrow E = \frac{1}{4\pi\epsilon_0} \cdot \frac{Q}{r^2}$$

$$= k_e \frac{Q}{r^2} \quad (r > R) \quad \propto \frac{1}{r^2}$$

(2) A point inside the sphere: Pick a gaussian surface as shown



$$\Phi_E = \oint \vec{E} \cdot d\vec{A} = \frac{q_{in}}{\epsilon_0}$$

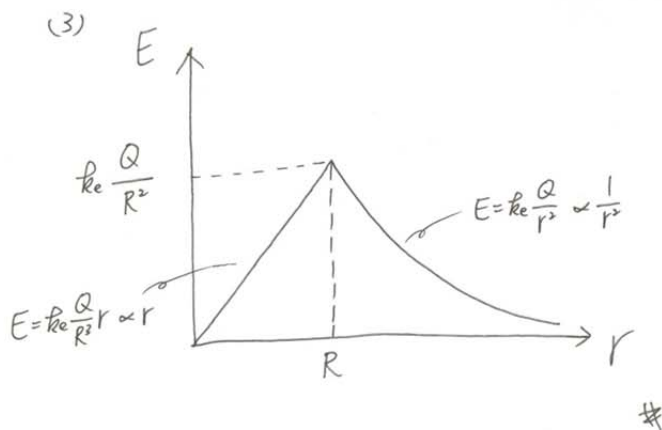
$$\Rightarrow E \cdot 4\pi r^2 = \frac{q_{in}}{\epsilon_0}, \quad q_{in} = P \cdot V$$

$$= \frac{Q}{\frac{4}{3}\pi R^3} \cdot \frac{4}{3}\pi r^3$$

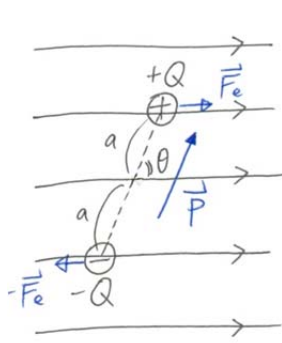
$$= Q \frac{r^3}{R^3}$$

$$\therefore E = \frac{1}{4\pi\epsilon_0} \frac{Q}{R^3} r$$

$$= k_e \frac{Q}{R^3} r \quad (r < R) \quad \propto r$$



3.



均匀电场 \vec{E}

$$P \equiv (2a) Q$$

- (1) $\left\{ \begin{array}{l} \text{Net force is zero.} \\ \text{Net torque is not zero. (The two forces produce a net torque on this dipole.)} \end{array} \right.$

$$\Rightarrow \boxed{\vec{\tau} \equiv \vec{r} \times \vec{F}}$$

$$\Rightarrow \tau = 2Fa \sin\theta$$

$$= \underset{\substack{\uparrow \\ F=QE}}{2QE} a \sin\theta = \underset{\substack{\uparrow \\ P=2aQ}}{PE} \sin\theta$$

$$\Rightarrow \vec{\tau} = \vec{P} \times \vec{E} \quad \#$$

(2) The work done to rotate the dipole is stored as potential energy in the system.

$$dW = \tau d\theta$$

$$U_f - U_i = \int_{\theta_i}^{\theta_f} \tau d\theta = \int_{\theta_i}^{\theta_f} PE \sin\theta d\theta = PE \underbrace{\int_{\theta_i}^{\theta_f} \sin\theta d\theta}_{-\cos\theta \Big|_{\theta_i}^{\theta_f}} = PE (\cos\theta_i - \cos\theta_f) \quad \#$$

(若令 $U_i = 0$ at $\theta_i = 90^\circ$, 则 $\cos\theta_i = 0$ & $U = U_f = -\vec{P} \cdot \vec{E}$)

4. |

4. The isolated system always tend toward disorder.

① Entropy is a measurement of the degree of the disorder of a system.

It is defined by

$$\Delta S = \int \frac{dQ}{T}$$

② For a adiabatic free expansion of a gas, if gas initial volume V_i expands to V_f . Then the ^{changes in} entropy can be defined as

$$\Delta S = \int_i^f \frac{dQ_r}{T} = \frac{1}{T} \int_i^f dQ_r$$

Since T is constant for this process, so we can't take $Q=0$ (even for adiabatic case $Q=0$), so we must choose an isothermal, reversible expansion here. In isothermal process

$$dQ = -dW = -P dV$$

Therefore,

$$\Delta S = - \int_{V_i}^{V_f} \frac{P}{T} dV = \frac{P}{T} \int_{V_i}^{V_f} dV$$

Let V_f is large than V_i

so $\Delta S = \text{positive}$

$$= \frac{nR}{V} \int_{V_i}^{V_f} dV = nR \ln \left(\frac{V_f}{V_i} \right)$$

For an ideal gas

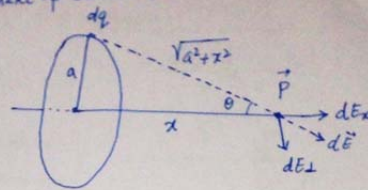
$$PV = nRT$$

$$\Delta S = nR \ln \left(\frac{V_f}{V_i} \right)$$

5.1

Q.5 (1) Let the direction of electric field is in x-direction -
 For the symmetry the perpendicular component of electric field will be canceled each other.

The parallel component of an electric field contribution from a segment of charge 'dq' on the ring



$$dE_x = k_e \frac{dq}{r^2} \cos\theta \quad \dots \text{where} \quad \textcircled{1}$$

$$= \frac{dq}{\sqrt{a^2+x^2}} \quad k_e = \frac{1}{4\pi\epsilon_0}$$

From the geometry, we can write

$$\cos\theta = \frac{x}{\sqrt{a^2+x^2}} \quad \dots \textcircled{2}$$

From eq. ① and ② we get

$$dE_x = k_e \frac{x dq}{(a^2+x^2)\sqrt{a^2+x^2}} = \frac{k_e x}{(a^2+x^2)^{3/2}} dq$$

So, from all segments of the ring, the electric field at P is

$$E_x = \int \frac{k_e x}{(a^2+x^2)^{3/2}} dq \quad \text{or} \quad \boxed{E = \frac{k_e x}{(a^2+x^2)^{3/2}} Q}$$

(2) If $x \gg a$, then the ring will act as a point charge so we can write $a=0$,

$$E = \frac{k_e x Q}{x^3} = k_e \frac{Q}{x^2}$$

$$\boxed{E = k_e \frac{Q}{x^2}}$$